


Geometric Quantization and Rep theory

- Main thm: Lie group G (semisimple or nilpotent)

\hookrightarrow adjoint rep $G \curvearrowright \mathfrak{g} \rightarrow$ Lie alg
 \hookrightarrow co-adjoint $G \curvearrowright \mathfrak{g}^*$

\hookrightarrow Irreducible rep of $G \longleftrightarrow$ — orbits of co-adjoint action
Thm 7.1.1 in the notes

by Kirillov.

character $\stackrel{\cong}{=} \int$ Fourier transform over the corresponding orbit.

- The theorem is an application of geometric quantization.

Q: Why this bijection is nice?

⊗ Describe explicitly the above bijection for $G = SO_2(\mathbb{R})$ (§ 1 of Venkatesh's notes)

⊗ Irreducible reps of $SO_3(\mathbb{R})$

- $SO_3(\mathbb{R}) = \left\{ A \in GL_3(\mathbb{R}) \text{ s.t. } \begin{array}{l} AA^T = A^T A \\ \det(A) = 1 \end{array} \right\} = \mathbb{I}$
 \searrow describe all rotations in \mathbb{R}^3 .

- $\mathcal{P}_n = \mathbb{R}$ -vector space of homogeneous
 polys of deg n in 3 variables
 e.g. $x^2y + yz^2, xyz \in \mathcal{P}_3$.

SO_3 $\subset GL_3 \curvearrowright \mathcal{P}_n$ by acting on x, y, z

$$\hookrightarrow g(x^2 + y^2 + z^2) = x^2 + y^2 + z^2 = r^2$$

$\forall g \in SO_3$

\Rightarrow have $SO_3(\mathbb{R})$ -equivariant map
 $xr^2: \mathcal{P}_{n-2} \rightarrow \mathcal{P}_n$

- $SO_3(\mathbb{R})$ -equivariant map $P_n \rightarrow P_{n-2}$
 by applying the Laplacian $\Delta = \partial_x^2 + \partial_y^2 + \partial_z^2$
 (by chain rule, $f \in P_n, g \in SO_3$
 want to show $\Delta(f \circ g) = g \circ \Delta(f)$
 $= \Delta g \circ \Delta f$)

- Interestingly, $P = \bigoplus P_n$

then xr^2 and Δ define a Lie alg
 action of \mathfrak{sl}_2 on P . \rightsquigarrow actually a

$\mathfrak{h} = [xr^2, \Delta] \dots$ more general fact
 will be discussed later.

- All ineps of $SO_3(\mathbb{R})$ are

$$V_n := \ker(\Delta: P_n \rightarrow P_{n-2})$$

$$\downarrow \cong P_n / r^2 P_{n-2}$$

$$\hookrightarrow P_n = V_n \oplus r^2 P_{n-2}$$

\hookrightarrow dropbox comment...

Fact: V_n 's are all the irreducible rep of $S^1_3(\mathbb{R})$, called spherical harmonics

eg: $V_2 = \text{span}_{\mathbb{R}}(xy, xz, yz, x^2 - y^2, y^2 - z^2)$.
 $(\partial_x^2 + \partial_y^2 + \partial_z^2)(x^2 - y^2) = 2 - 2 = 0$.

Rmk: V_n are eigenspaces for the Riemannian Laplacian $\approx L^2(S^2) =$ square-int func on S^2
 \hookrightarrow roughly Δ restricts on S^2 .
(drop box comment).

⊗ Characters of V_n 's

- Recall: $G \xrightarrow{P} GL(V)$ finite dim then
character $\text{Ch}: g \mapsto \text{tr}(P(g))$.

- $P_n = V_n \oplus \mathbb{R}^2 P_{n-2}$, since trace is

additive so $\text{Tr}_{P_n}(g) = \text{Tr}_{V_n}(g) + \text{Tr}_{\mathbb{R}^2 P_{n-2}}(g)$

$\Rightarrow \text{Ch}_{V_n} = \text{Ch}_{P_n} - \text{Ch}_{P_{n-2}}$ $\text{Tr}_{P_{n-2}}(g)$.

- Any $g \in SO_3(\mathbb{R})$ is conjugate to a rotation around z -axis at angle θ

$$g_\theta = \begin{matrix} x & y & z \\ \begin{pmatrix} \cos \theta & -\sin \theta & 0 \\ \sin \theta & \cos \theta & 0 \\ 0 & 0 & 1 \end{pmatrix} \end{matrix}$$



Claim: $\text{Ch}_{V_n}(g_\theta) = e^{in\theta} + e^{i(n-1)\theta} + \dots + e^{-in\theta}$

Check for V_2 : $\text{ch}_{P_2} = P_2 = \{x^2, y^2, z^2, xy\}$

Compute trace

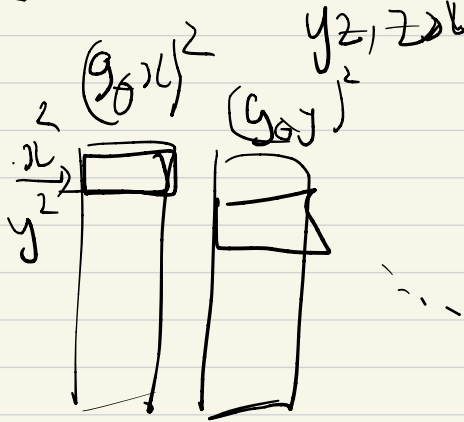
$$[x^2] (g_\theta x)^2$$

//

$$[x^2] (x \cos \theta + y \sin \theta)^2 = \cos^2 \theta$$

$$[y^2] (g_\theta y)^2 = \cos^2 \theta$$

$$[z^2] (g_\theta z)^2 = 1$$



$$\begin{aligned} & [xy] (g_\theta x)(g_\theta y) \\ &= [xy] (x \cos \theta + y \sin \theta)(-x \sin \theta + y \cos \theta) \\ &= \cos^2 \theta - \sin^2 \theta = 2 \cos^2 \theta - 1 \end{aligned}$$

$$\bullet [yz](g_b)(g_b z) = \cos \theta$$

$$[z1] \circ \stackrel{z}{=} \cos \theta$$

$$\text{ch}_{P_2}(g_b) = 4 \cos^2 \theta + 2 \cos \theta$$

$$\text{ch}_{P_0}(g_b) = 1$$

$\stackrel{IR}{=}$

$$\begin{aligned} \hookrightarrow \text{ch}_{V_2}(g_b) &= 4 \cos^2 \theta + 2 \cos \theta - 1 \\ &= 2(\cos 2\theta + 1) + 2 \cos \theta - 1 \\ &= 2 \cos 2\theta + 2 \cos \theta + 1 \\ &= e^{2i\theta} + e^{-2i\theta} + e^{i\theta} + e^{-i\theta} + 1 \end{aligned}$$

$$\begin{aligned} - \chi_n(g_b) &= e^{ni\theta} + e^{i(n-1)\theta} + \dots + e^{-ni\theta} \\ &= \frac{e^{i(n+1)\theta} - e^{-i(n+1)\theta}}{e^{i\theta} - e^{-i\theta}} \end{aligned}$$

$$= \frac{(e^{i(n+1/2)\theta} - e^{-i(n+1/2)\theta})}{(e^{i\theta/2} - e^{-i\theta/2})}$$

⊕ Show χ_n 's are all irreps of $SO_3(\mathbb{R})$.
 by using $\chi_n \dots \rightarrow L^2(G)$

- Recall: G finite group $\rightarrow G$ -invariant inner $\langle \cdot, \cdot \rangle$
 product on $\text{Fun}(G, \mathbb{C}) \supset$ Class functions

Irre characters form an orthonormal basis
 wrt $\langle \cdot, \cdot \rangle$ on space of class functions

$$\chi \text{ irred} \Leftrightarrow \langle \chi, \chi \rangle = 1$$

- G -invariant inner product $\langle \cdot, \cdot \rangle$ on $L^2(G)$
 comes from Haar measure μ on G

$$\text{i.e. } \langle f, g \rangle = \int_G f(x) \overline{g(x)} d\mu(x)$$

= What is μ for $SO_3(\mathbb{R})$?

• Recall Riemann's class: Define bi-invariant
 metric g on $T_e SO_3(\mathbb{R}) = \mathfrak{so}_3$.

\hookrightarrow Killing form on \mathfrak{so}_3 . $\left\{ \frac{\partial}{\partial J_1}, \dots, \frac{\partial}{\partial J_2} \right\}$

Find 3-dim G -invariant diff form

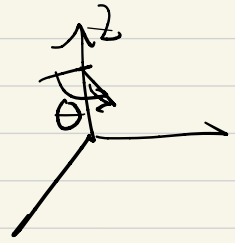
$$\sqrt{\det(g_{ij})} dJ_x \wedge dJ_y \wedge dJ_z$$

* Local coordinates on $SO_3(\mathbb{R})$;

every element in $SO_3(\mathbb{R})$ corresponds
to a rotation, which is uniquely determined
by a rotation axis and a rotation angle α

use
spherical
coordinates (θ, φ)

S^2



To represent a matrix in $SO_3(\mathbb{R})$ using
 $(\theta, \varphi, \alpha) \rightsquigarrow$ use exponential map
 $S^3 \rightarrow SO_3(\mathbb{R})$

To do: Write down ω explicitly ...

Check : χ_n 's are orthonormal

$\Rightarrow V_n$'s are all irreducible reps of $SO_3(\mathbb{R})$

- Need to show χ_n 's is dense

In the space of class functions
 $\subset L^2(SO_3(\mathbb{R}))$. □