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# Talk 8 - §2 of Venkatesh's notes

⊗ Recall:  $V_N$  irrep of  $SO_3(\mathbb{R})$  of dim  $2N+1$   
 $\chi_N$  char of  $V_N$

Then  $(\chi_N \sqrt{J})(e^X) = \frac{e^{i(N+\frac{1}{2})s} - e^{-i(N+\frac{1}{2})s}}{is}$

where  $X \in \mathfrak{so}_3$ ,  $|X| = s \xrightarrow{\text{wrt } J_x, J_y, J_z \text{ basis}}$

$J$  Jacobian at  $X$  of  $\exp: \mathfrak{so}_3 \rightarrow SO_3(\mathbb{R})$

↳ need to check ...

• Denote  $\hat{\mu}: \mathbb{R}^3 \rightarrow \mathbb{R}$  be the Fourier transform of the area measure of  $S^2_R \subset \mathbb{R}^3$ .

then  $\hat{\mu}(k) = \frac{e^{iR|k|} - e^{-iR|k|}}{i|k|} \cdot 2\pi R$   
 $\lambda = |k|, k \in \mathbb{R}^3$

⇒ If we normalise  $\mu$  s.t. factor  $2\pi R$  disappears and the total area is  $2N+1 = \dim V_N$  then

(\*)  $(\chi_N \sqrt{J})(e^X) = \int_{S^2_{N+\frac{1}{2}}} e^{i\langle \xi, X \rangle} d\xi$

↳  $\underline{Q}$ : Where does  $S^2_{N+\frac{1}{2}}$  come from?

Today: Ignore jacobian  $j$ , focus on explaining  $S_{N+1/2}^2$  and  $\mu \dots$

Corollary:  $X=0$

$$Z_{N+1} = \text{dim } V_N \stackrel{X=0}{=} (K_N \sqrt{j})(e^0) = \int_{S_{N+1/2}^2} d\zeta = \text{area } S_{N+1/2}^2$$

### Speculations of (\*)

① From corollary, there exists  $\{v_1, \dots, v_N\}$  basis of  $V_N$  in bijection with partition

$$S_{N+1/2}^2 = \bigsqcup_{k=1}^{2N+1} \mathcal{O}_k$$

with  $\mu(\mathcal{O}_k) = 1$ .

② If each  $v_k$  satisfies  $\zeta_k \in \mathcal{O}_k$   
 $S_{\mathcal{O}_3} \otimes V_N \xrightarrow{\text{diff}}_{S_{\mathcal{O}_3} \otimes V_N} X \cdot \zeta_k \approx i \langle \zeta_k, X \rangle v_k$  for some  
 then  $e^X \cdot \zeta_k \approx e^{i \langle \zeta_k, X \rangle} \zeta_k$

$$e^{\begin{bmatrix} v_1 & & \\ & \ddots & \\ & & v_k \end{bmatrix}} = e^{\begin{bmatrix} i \langle \xi_1, X \rangle & & \\ & \ddots & \\ & & i \langle \xi_k, X \rangle \end{bmatrix}}$$

then  $\text{Tr}(e^X) = \sum e^{i \langle \xi_k, X \rangle} = \int_{S_{N+\frac{1}{2}}^2} e^{\langle \xi, X \rangle} d\xi$

$\Rightarrow$  Decomposition  $S_{N+\frac{1}{2}}^2 = \bigsqcup \mathcal{O}_k$

correspond to diagonalisation of  $SO_3(\mathbb{R})$ -action

- But the above picture is not correct

as  $v_N$  is then sum of 1-dim rep

- Better way to speculate:  $v_k$  can be

decomposed as sum of  $X$ -eigenvectors,  
 each with eigenvalue  $i \langle \xi, X \rangle$  for some  $\xi \in \mathcal{O}_k$ .

Why would this speculation  $\rightarrow [X, Y] \cdot v_k = i \langle \xi_{(v)} [X, Y] \rangle v_k \neq 0$

speculation  $[X, Y] v_k = X Y v_k - Y X v_k = 0$

fix this issue?

$$i_X : \mathcal{O}_k \rightarrow \mathbb{R} \quad \xi \mapsto \langle \xi, X \rangle$$

$\uparrow$   
 $\ell_X$

$$v_k = \sum_{i=1}^r \alpha_i t_i \quad \curvearrowright \quad X \cdot t_i = \langle \xi_i, X \rangle t_i$$

⊗ General formula :

-  $G$  Lie grp,  $\xi \in \mathfrak{g}^*$   $\text{ad} \cdot g \rightarrow \text{---} [ ]$

Adjoint action  $\text{Ad} : G \rightarrow \text{GL}(\mathfrak{g})$

- Co-adjoint action  $\text{Ad}^* : G \rightarrow \text{GL}(\mathfrak{g}^*)$

by  $\langle \text{Ad}(g)\vartheta, f \rangle = \langle \vartheta, \text{Ad}^*(g^{-1})f \rangle$

- Let  $\mathcal{O}$  be orbit of  $\xi \in \mathfrak{g}^*$ .  
 $\hookrightarrow$  manifold (?)

$\downarrow$  differentiate

$$\text{ad}^* : \mathfrak{g} \rightarrow \text{GL}(\mathfrak{g}^*)$$

$$\gamma : X \mapsto \vartheta \in \mathfrak{g}^*$$

$$\gamma'(0) = X, \gamma(0) = 1$$

$$\leftarrow \exp(tX) : \mathbb{I} \rightarrow G$$

$$\text{ad}^*(X)\vartheta := \left. \frac{d}{dt} \right|_{t=0} \text{Ad}^*(\exp(tX))\vartheta$$

- Claim :  $\text{Tan}_\xi \mathcal{O} = \{ \text{ad}^*_X(\xi) : X \in \mathfrak{g} \}$

Proof : path  $\gamma_X(t) = \text{Ad}^*_{\exp(tX)}(\xi)$   $X \in \mathfrak{g}$   
 $\hookrightarrow$  exp of  $X$

$$\gamma(0) = \text{Ad}_1^*(\xi) = \xi.$$

$$\gamma'_X(0) = \text{ad}_X^*(\xi) \in T_\xi(\mathcal{O}).$$

Thus, we have defined a linear map

$$\mathfrak{g} \rightarrow \text{Tan}_\xi(\mathcal{O})$$

$$X \mapsto \gamma'_X(0)$$

Its kernel <sub>a</sub> those  $X \in \mathfrak{g}$  s.t.  $\text{ad}_X^*(\xi) = 0$ .

$$\text{ad}_{[X,Y]}^* = [\text{ad}_X^*, \text{ad}_Y^*] \quad \mathfrak{g}_\xi \text{ Lie sub alg of } \mathfrak{g}.$$

-  $\mathfrak{g}_\xi$  is Lie alg of  $G_\xi$ , stabiliser of  $\xi$  under  $\text{Ad}^*$ .

Indeed,  $\exp(\mathfrak{g}_\xi) \subset G_\xi \rightarrow T_{\text{id}} G_\xi \supset \mathfrak{g}_\xi$ .

Conversely, if  $\gamma: I \rightarrow G_\xi \rightarrow \gamma(t) \in G_\xi$

$$\text{i.e. } \text{Ad}_{\gamma(t)}^* \xi = \xi \quad \forall t.$$

$$\frac{d}{dt} \Big|_{t=0} \left( \text{ad}_{\gamma(t)}^* \xi = \xi \right) = 0 \Rightarrow \gamma'(0) \in \mathfrak{g}_\xi.$$

- Injection  $\mathfrak{g}/\mathfrak{g}_\xi \hookrightarrow \text{Tan}_\xi(\mathcal{O})$

$\mathcal{O} = G/G_\xi$  so  $\mathfrak{g}/\mathfrak{g}_\xi$  and  $\text{Tan}_\xi(\mathcal{O})$  have same dim.  $\Rightarrow$  surjective ▣

$$\mathcal{O} = G \cdot \xi$$

— There is a natural  $G$ -invariant non-deg symplectic form  $\omega$  on  $\mathcal{O}$  (i.e. an alternating non-deg bilinear form on  $T_x\mathcal{O}$ ) defined by

For each  $t_1, t_2 \in T_{g \cdot \xi} \mathcal{O}$ , choose  $X, Y \in \mathfrak{g}$  s.t.  $X \cdot \xi = t_1, Y \cdot \xi = t_2$ .

$$\omega(t_1, t_2) := \int ([X, Y]) \cdot \xi \quad \text{ad}^*(X) \cdot \xi$$

well-defined, does not depend on the choice of  $X$  and  $Y$ .

canonical, i.e. does not depend on  $\xi$

replace  $\xi$  by anything in  $\mathcal{O}$ .

—  $\omega$  is  $G$ -invariant:  $g \in G$ , on  $T_{g \cdot \xi}(\mathcal{O})$

$$\text{we have } \omega_{g \cdot \xi}(g_* t_1, g_* t_2) = \omega_{\xi}(t_1, t_2)$$

where  $g_* : T_{\xi}(\mathcal{O}) \rightarrow T_{g \cdot \xi}(\mathcal{O})$  because  $g : \mathcal{O} \rightarrow \mathcal{O}$   
 $\downarrow \text{ad}^*$   $\downarrow \text{Ad}^*$

$\Rightarrow (\mathcal{O}, \omega)$  is a symplectic manifold.

• Example: - Adjoint action of  $SO_3$  on  $\mathfrak{so}_3$  is iso to the rotation action of  $SO_3$  on  $\mathbb{R}^3$

$\Rightarrow$  Co-adjoint action (conjugate transpose of adjoint action) is also rotation on  $\mathbb{R}^3$

$\Rightarrow$  Co-adjoint orbits correspond to Spheres in  $\mathbb{R}^3$   
 $\circlearrowleft$  always  $\dim \cong 2$   
 even dim.

Theorem (Kirillov, others)  $G$  connected Lie grp (either nilpotent or semisimple).  $\pi$  tempered rep of  $G$  (i.e.  $\pi$  lies in  $L^2(G) \rtimes G$ ). Then there is an orbit  $\mathcal{O}$  of  $G$  on  $\mathfrak{g}^*$  s.t.  $(X + \mathfrak{h}) \cdot (e^X) =$  Fourier trans of  $\left(\frac{\omega}{2\pi}\right)^d$  on  $\mathcal{O}$   $\omega$  symplectic form to  $\mathcal{O}$   $2d = \text{real dim of } \mathcal{O}$ .

Upshot: } Know where  $\mathcal{O}$  comes from  
 } — how to define the measure

$\otimes$  Describe — explicitly for  $SO_3$ :

(exer 2.3.7, p. 8)

next time





